

# Study of the couplings of QED and QCD from the Adler function

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Harvey Meyer and Hartmut Wittig

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# $\bar{a}_\mu^{\text{HLO}}(Q_{\text{ref}}^2)$

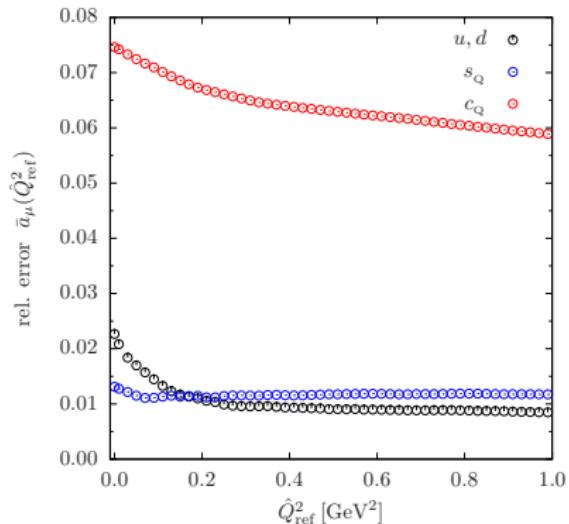
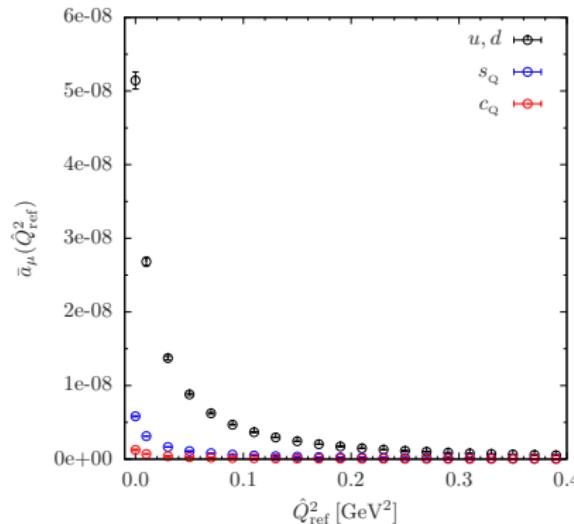
[talk by Hanno Horch]

$\bar{a}_\mu^{\text{HLO}}$  is dominated by the low  $Q^2$  region : noisy and long-distance contributions

$$\bar{a}_\mu^{\text{HLO}}(Q_{\text{ref}}^2) \equiv 4\pi^2 \left(\frac{\alpha}{\pi}\right)^2 \int_{Q_{\text{ref}}^2}^{\infty} dQ^2 f(Q^2, m_\mu^2) [\Pi(Q^2) - \Pi(Q_{\text{ref}}^2)] \xrightarrow{Q_{\text{ref}}^2 \rightarrow 0} a_\mu^{\text{HLO}}$$

integrand is peaked at  $Q^2 \sim m_\mu^2$

$m_\mu^2 \sim 0.01 \text{ GeV}^2$



~ here we will consider physical quantities sensitive to larger  $Q^2$  regime

# running of QED coupling

$$\Delta\alpha_{\text{QED}}^{\text{had}}$$

# $\Delta\alpha_{\text{QED}}^{\text{had}}$

$$\alpha(Q^2) = \frac{\alpha}{1 - \Delta\alpha_{\text{QED}}(Q^2)}$$

- ▶ vacuum polarisation: charge screening  $\rightsquigarrow$  running of QED coupling
- ▶ Standard Model (SM) precision tests and sensitivity to new physics requires precise knowledge of  $\Delta\alpha_{\text{QED}}(Q^2)$ : input parameter of SM
- ▶  $\alpha = 1/137.035999074(44)$  [0.3 ppb] [PDG, 2013]
- ▶  $\alpha(M_Z^2) = 1/128.952(14)$   $[10^{-4}] \rightsquigarrow 10^5$  less accurate ... [M. Davier et al., 1010.4180]
- ▶ uncertainty in  $\alpha(M_Z^2)$  is significantly larger than that of  $M_Z$
- ▶ hadronic effects:  $\alpha(Q^2)$  depends strongly on  $Q^2$  at low energies  
hadronic uncertainties propagate ...



$$\Delta\alpha_{\text{QED}}^{\text{had}}$$

$$\alpha(Q^2) = \frac{\alpha}{1 - \Delta\alpha_{\text{QED}}(Q^2)}$$

- ▶ leading order (LO) contribution



$$\int d^4x e^{iQx} \langle J_\mu(x) J_\nu(0) \rangle = (Q_\mu Q_\nu - Q^2 \delta_{\mu\nu}) \Pi(Q^2)$$

$$J_\mu(x) = \sum_{f=1}^{N_f} Q_f \bar{\psi}_f(x) \gamma_\mu \psi_f(x)$$
$$Q_f \in \{-1/3, 2/3\}$$

- ▶  $\Pi(Q^2)$ : photon **vacuum polarisation function (VPF)**

$$\Delta\alpha_{\text{QED}}^{\text{had}}$$

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- ▶  $\Pi(Q^2)$ : photon vacuum polarisation function (VPF)

$$\Delta\alpha_{\text{QED}}(Q^2) = 4\pi\alpha \left( \Pi(Q^2) - \Pi(0) \right)$$

- ▶ Adler function  $D(Q^2)$ :

$$\begin{aligned} \frac{D(Q^2)}{Q^2} &= 12\pi^2 \frac{d\Pi(q^2)}{dq^2} \\ &= -\frac{3\pi}{\alpha} \frac{d}{dq^2} \Delta\alpha_{\text{QED}}^{\text{had}}(q^2) \end{aligned} \quad Q^2 = -q^2$$

$$\Delta\alpha_{\text{QED}}^{\text{had}}$$

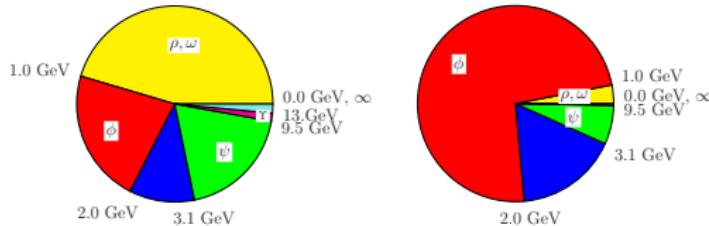
$$\alpha(Q^2) = \frac{\alpha}{1 - \Delta\alpha_{\text{QED}}(Q^2)}$$

- ▶ Combine experimental data and perturbation theory (PT)

$$\begin{aligned}\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) &= \Delta\alpha_{\text{had}}^{(5)}(-M_0^2)^{\text{exp}} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(-M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_0^2)]^{\text{pQCD}} \\ &+ [\Delta\alpha_{\text{had}}^{(5)}(M_Z^2) - \Delta\alpha_{\text{had}}^{(5)}(-M_Z^2)]^{\text{pQCD}}\end{aligned}$$

b)

$$M_0^2 = (2.5 \text{ GeV})^2 \approx 6 \text{ GeV}^2$$



$\gtrsim [1\%]$

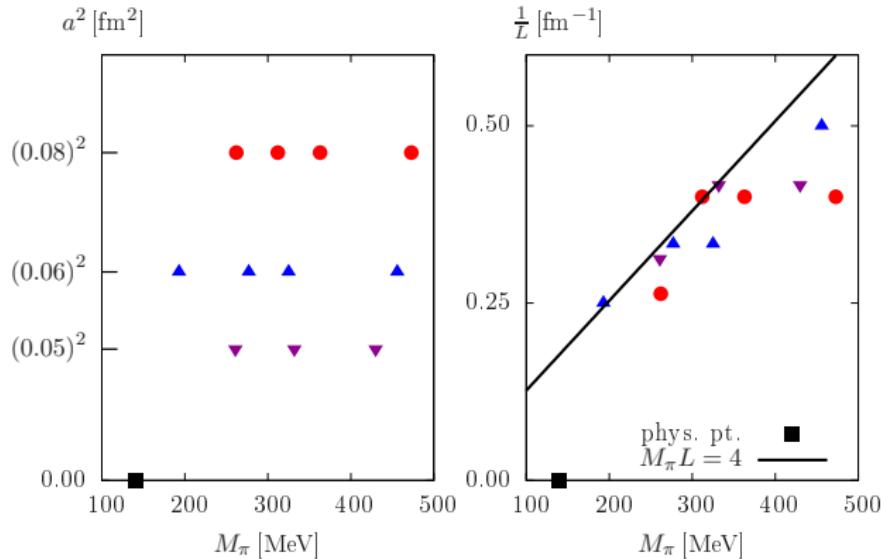
[F. Jegerlehner, 0807.4206]

- ▶ can lattice QCD reach similar precision for  $\Delta\alpha_{\text{QED}}^{\text{had}}$ ?

# lattice setup

► [Mainz, 1112.2894]

[talk by Hanno Horch]



$N_f = 2$   $\mathcal{O}(a)$  improved Wilson fermions [CLS]

increased statistics

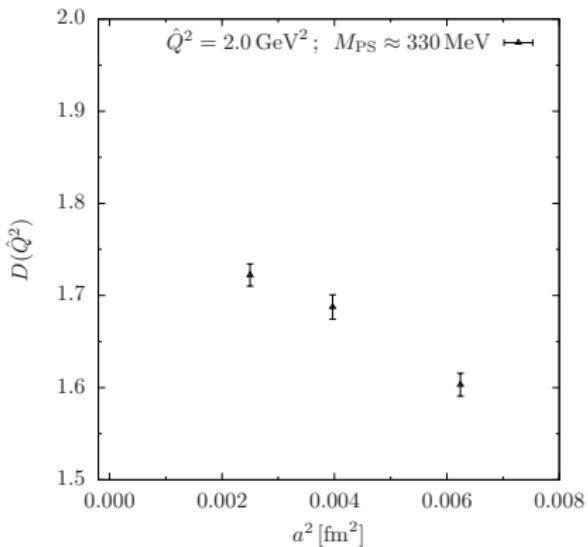
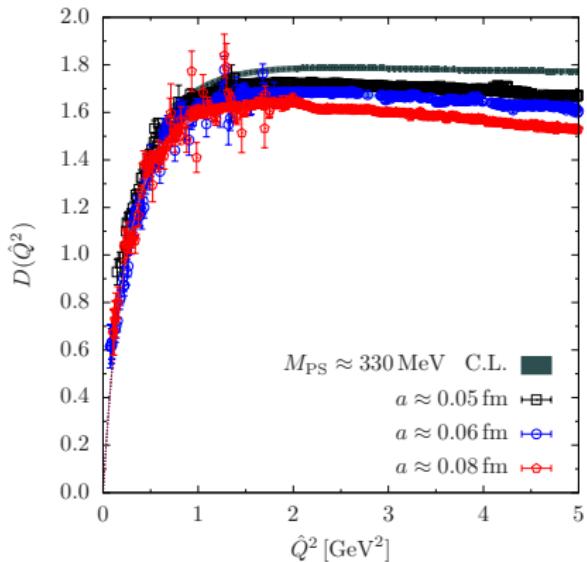
strange and charm are quenched :  $s_0, c_0$

only quark-connected contributions

# Adler function : lattice spacing dependence

$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$

*u, d*



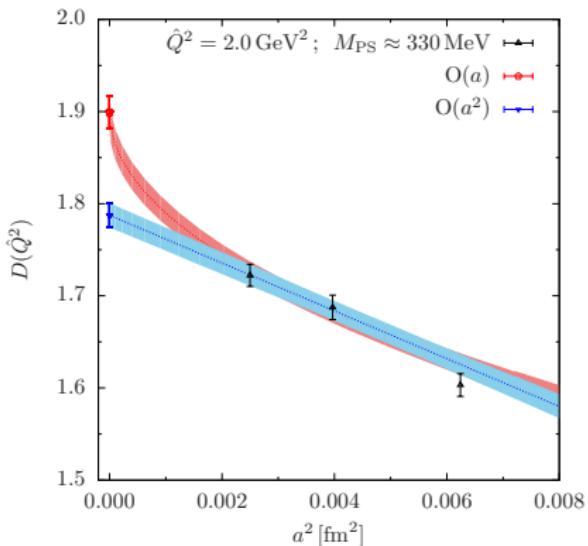
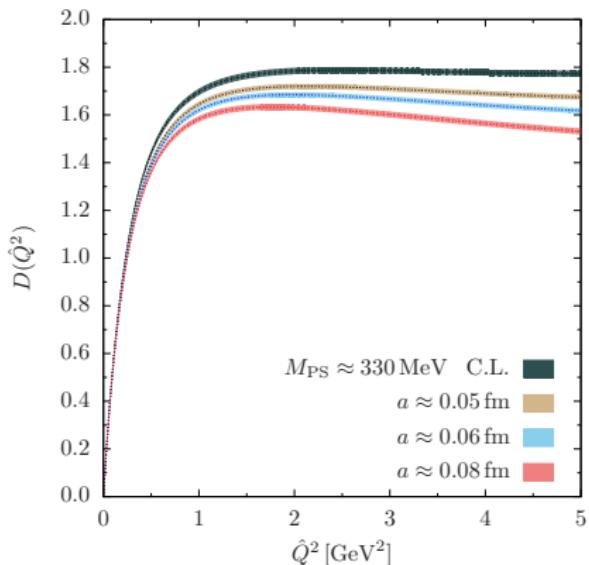
combined fits [\[talk by Hanno Horch\]](#)

lattice artefacts : dominant systematic effect for  $Q^2 \gtrsim 1 \text{ GeV}^2$

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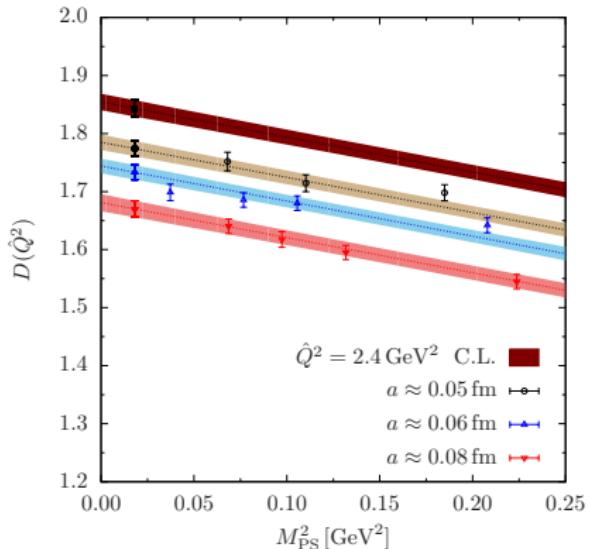
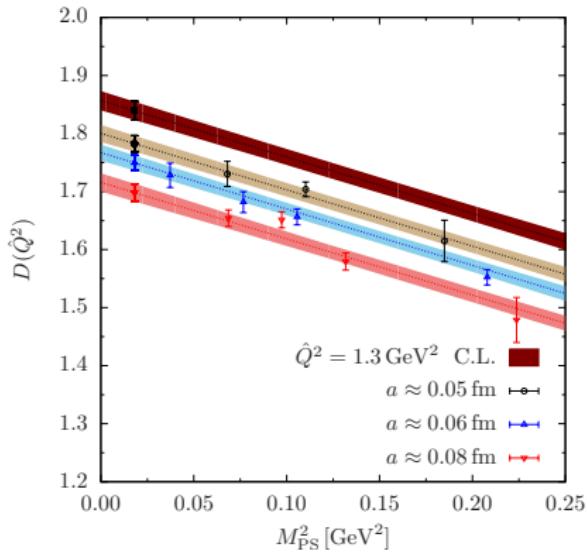
combined fits [\[talk by Hanno Horch\]](#)

lattice artefacts : dominant systematic effect for  $Q^2 \gtrsim 1 \text{ GeV}^2$

# Adler function : light-quark mass dependence

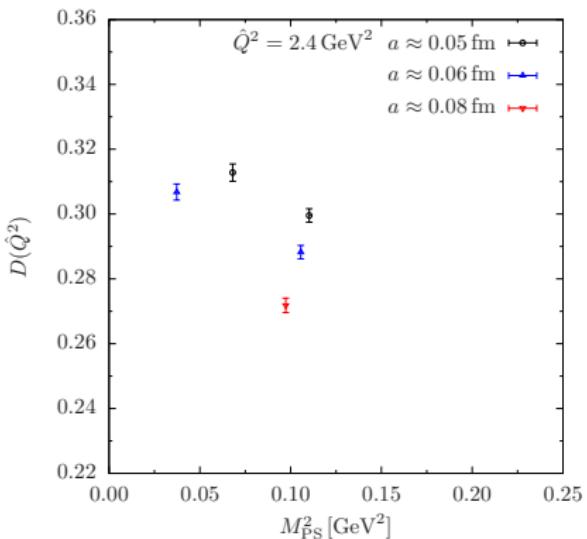
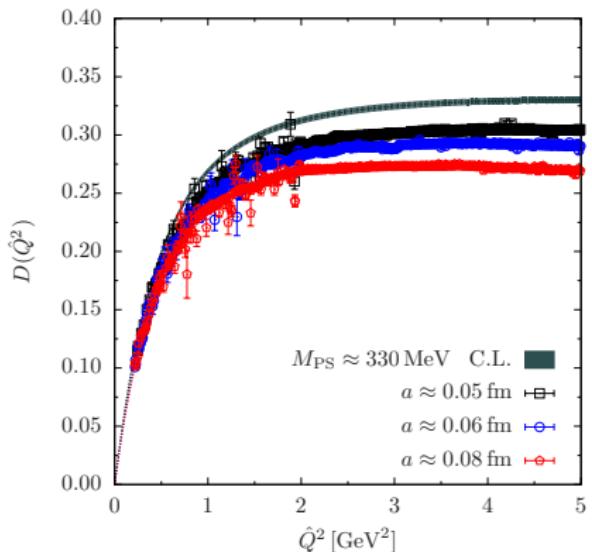
$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$

*u, d*



# Adler function : strange quark

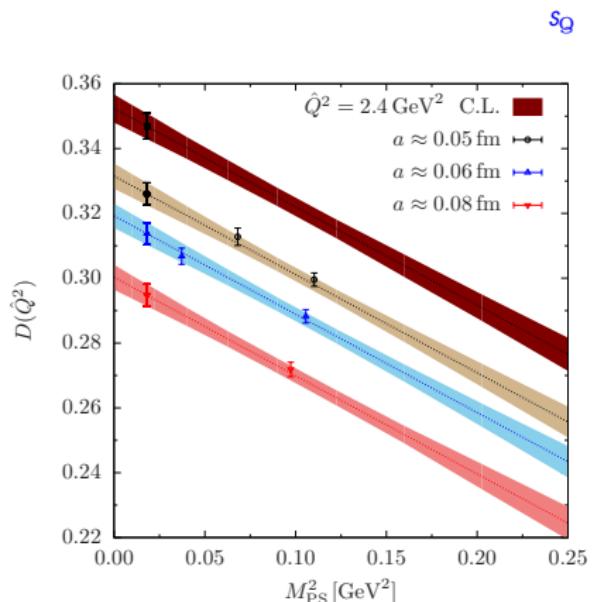
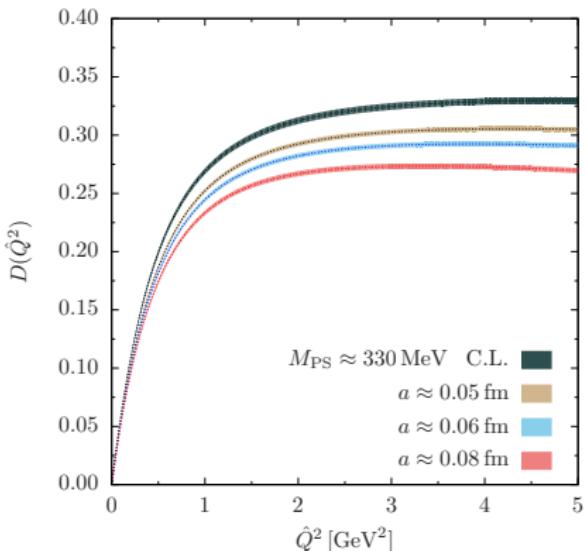
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lattice artefacts : dominant systematic effect for  $Q^2 \gtrsim 1$  GeV $^2$

# Adler function : strange quark

$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$



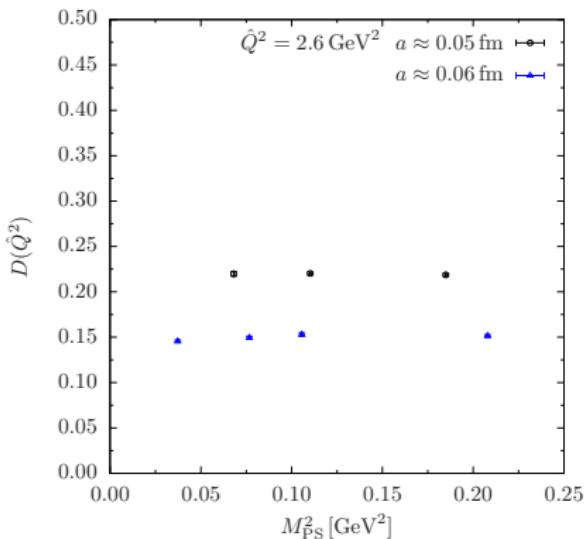
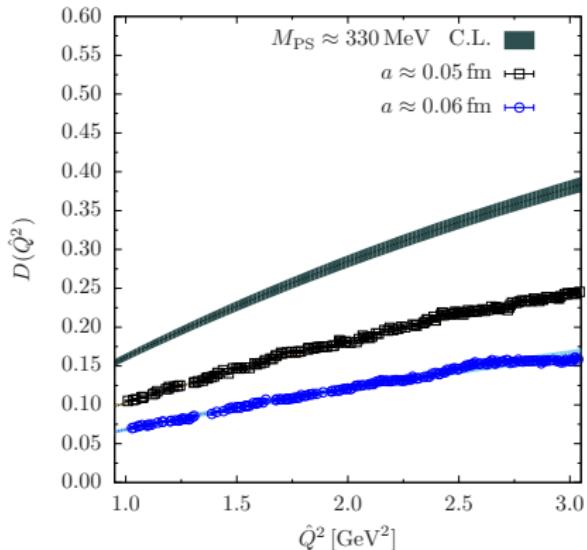
lattice artefacts : dominant systematic effect for  $Q^2 \gtrsim 1 \text{ GeV}^2$

[PRELIMINARY]

# Adler function : charm quark

$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$

$c_0$

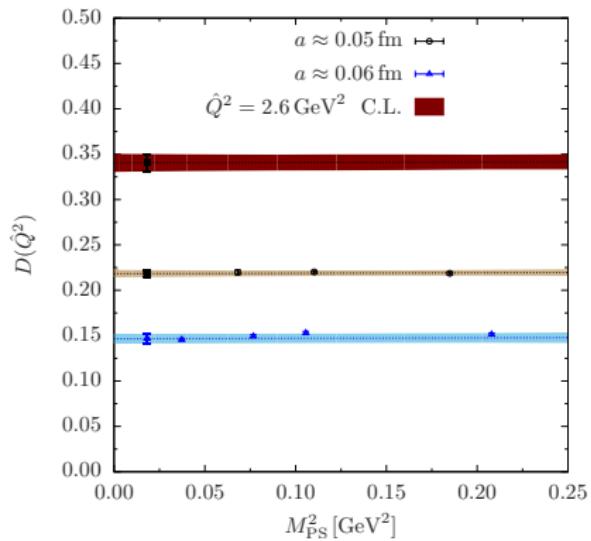
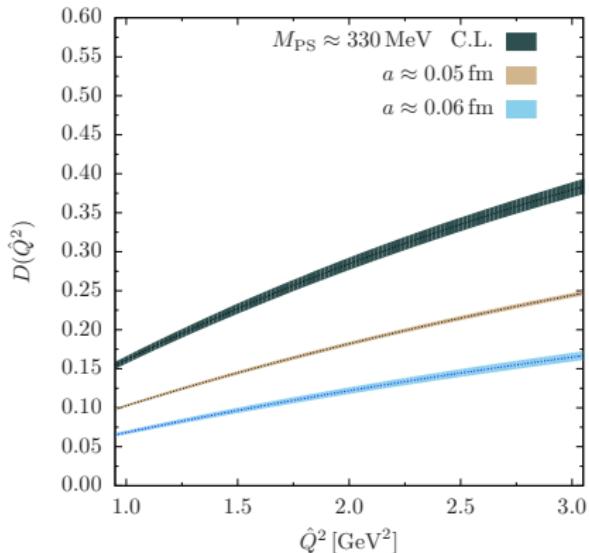


lattice artefacts : dominant systematic effect for  $Q^2 \gtrsim 1 \text{ GeV}^2$

# Adler function : charm quark

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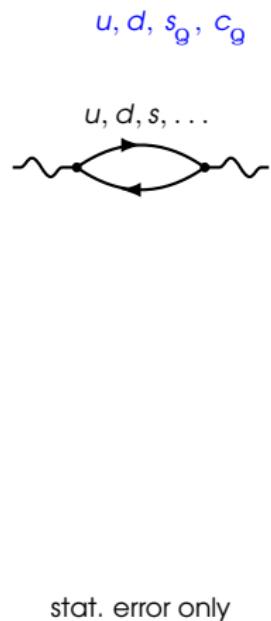
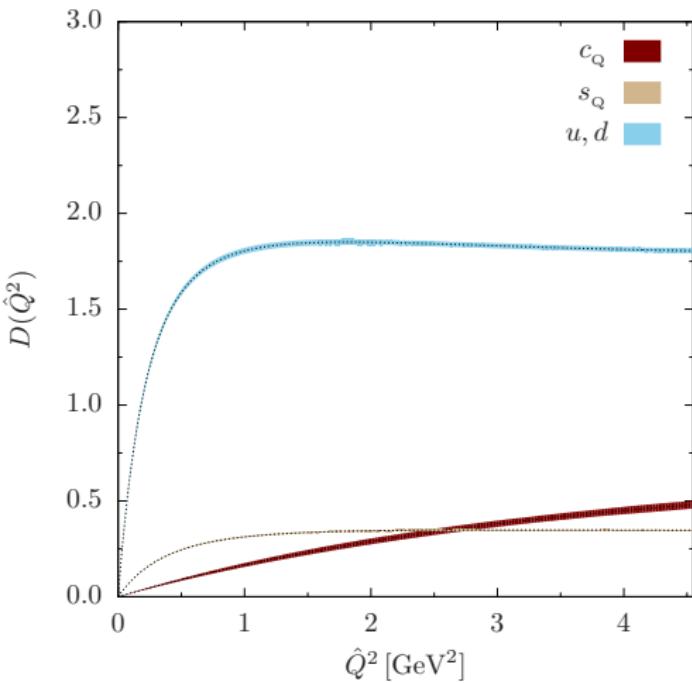
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[PRELIMINARY]

# Adler function : flavour contributions

$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$

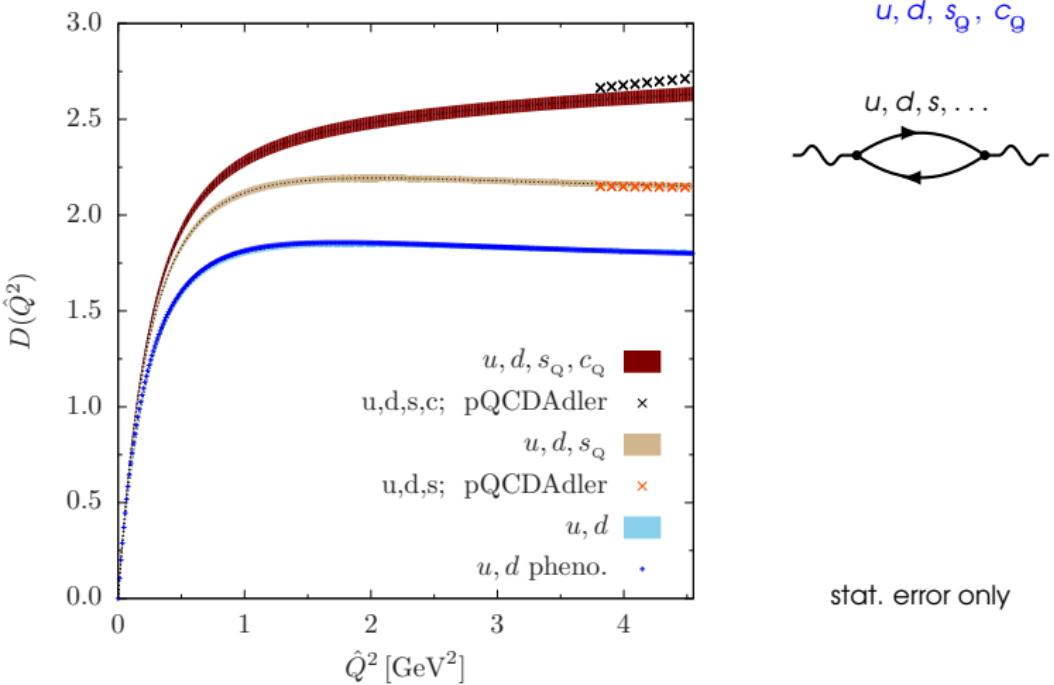
$\alpha \rightarrow 0$   
 $M_\pi^{\text{phys}}$



# Adler function : flavour contributions

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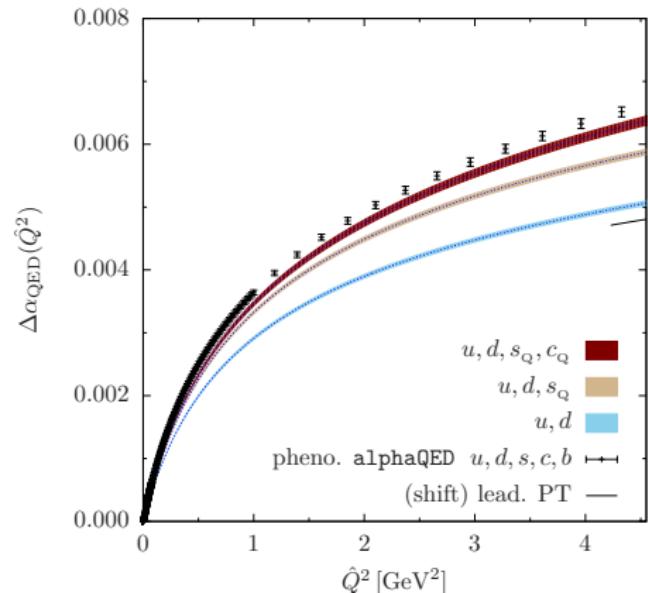
$\alpha \rightarrow 0$   
 $M_\pi^{\text{phys}}$



# running QED coupling: $\Delta\alpha_{\text{QED}}^{\text{had}}(Q^2)$

$$\alpha(Q^2) = \frac{\alpha}{1 - \Delta\alpha_{\text{QED}}(Q^2)}$$

$$\Delta\alpha_{\text{QED}}^{\text{had}}(Q^2) = 4\pi\alpha \left( \Pi(Q^2) - \Pi(0) \right)$$



$u, d$

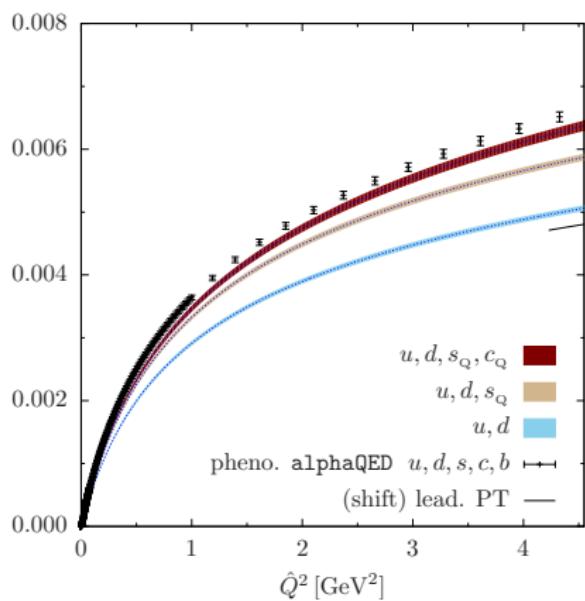
$u, d, s_Q$

$u, d, s_Q, c_Q$

$u, d, s, c, b$

# running QED coupling: $\Delta\alpha_{\text{QED}}^{\text{had}}(Q^2)$

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$$\Delta\alpha_{\text{QED}}^{\text{had}}(Q^2) = 4\pi\alpha \left( \Pi(Q^2) - \Pi(0) \right)$$

[PRELIMINARY]

►  $\Delta\alpha_{\text{QED}}^{\text{had}}(1 \text{ GeV}^2)$

[ $10^{-3}$  units]

$u, d :$  2.95(04)(05) [2%]

$u, d, s_Q :$  3.36(04)(06) [2%]

$u, d, s_Q, c_Q :$  3.50(05)(09) [3%]

$u, d, s, c, b :$  3.64(04) [1%]

Pheno. [alphaQED package, F. Jegerlehner]

► difference :  $\Delta\alpha_{\text{QED}}^{\text{had}}(4 \text{ GeV}^2) - \Delta\alpha_{\text{QED}}^{\text{had}}(1 \text{ GeV})$

$u, d :$  1.97(02)(06) [3%]

$u, d, s_Q :$  2.33(02)(08) [3%]

$u, d, s_Q, c_Q :$  2.65(02)(15) [6%]

$u, d$

$u, d, s_Q$

$u, d, s_Q, c_Q$

$u, d, s, c, b$

# comparison to perturbative QCD

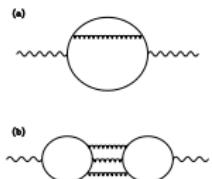
$$\alpha_s$$

# Operator Product Expansion (OPE)

Matching of lattice determinations of the VPF and Adler function to perturbation theory

Non-singlet and singlet contributions

$$D^{(N_f)}(\hat{Q}^2, \alpha_s) = \sum_f Q_f^2 D_{\text{con}}(\alpha_s, \hat{Q}^2, m_f) + \sum_{f,f'} Q_f Q_{f'} D_{\text{disc}}(\alpha_s, \hat{Q}^2, m_f, m_{f'})$$



OPE

$$\begin{aligned} D_{\text{con}}^{\text{OPE}}(\hat{Q}^2, \alpha_s, m_f) &= D_0(\alpha_s, \hat{Q}^2, \mu^2) \\ &+ D_2^m(\alpha_s, \hat{Q}^2, \mu^2) \frac{(m_f[\hat{Q}^2])^2}{\hat{Q}^2} \\ &+ D_4^F(\alpha_s, \hat{Q}^2, \mu^2) \frac{m_f \langle \bar{\psi}_f \psi_f \rangle}{\hat{Q}^4} \\ &+ D_4^G(\alpha_s, \hat{Q}^2, \mu^2) \frac{\langle O_{\text{OPE}}^{(4)} \rangle}{\hat{Q}^4} \\ &+ \mathcal{O}\left(\frac{1}{\hat{Q}^6}\right) \end{aligned}$$

# Operator Product Expansion (OPE)

- ▶ only quark-connected contributions **both** in PT and lattice
- ▶ Wilson coefficients  $D_0$ ,  $D_2^m$ ,  $D_4^F$ ,  $D_4^G$  are computed in PT

$$D_0: \mathcal{O}(\alpha_s^4), \quad D_2^m: \mathcal{O}(\alpha_s^2), \quad D_4^F: \mathcal{O}(\alpha_s^2), \quad D_4^G: \mathcal{O}(\alpha_s)$$

connection of  $\alpha_s$  to  $\Lambda_{\overline{\text{MS}}}^{(N_f=2)}$  via the 4-loop  $\beta$ -function

- ▶ fit of lattice data to PT: **range of validity of PT vs. discretisation effects**
- ▶ fit parameters:  $\alpha_s(\mu = 2 \text{ GeV})$ ,  $\langle O_{\text{OPE}}^{(4)} \rangle$   
and 2 parameters for lattice artefacts

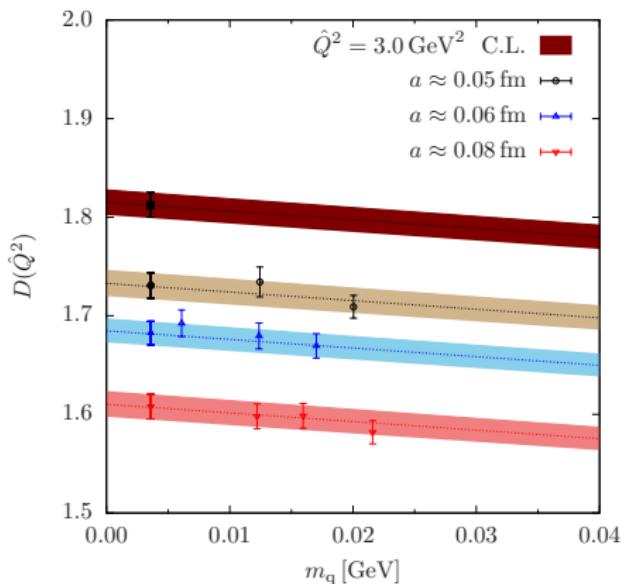
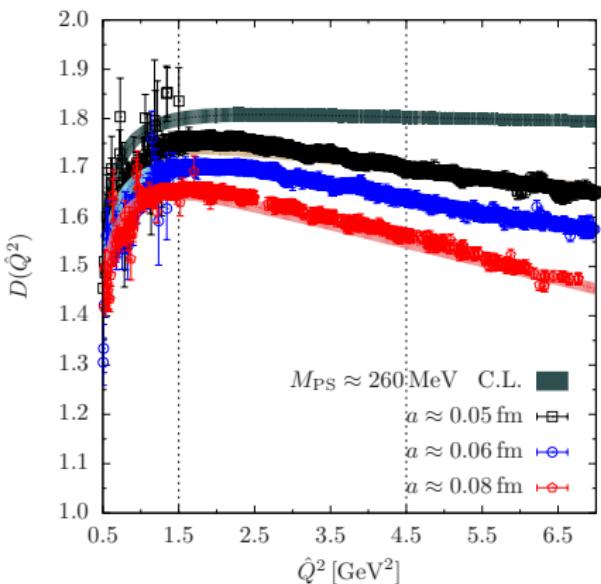
chiral condensate  $[\overline{\text{MS}}; \mu = 2 \text{ GeV}]$  from [\[FLAG. 1310.8555\]](#)

earlier lattice studies [\[JLQCD, 0807.0556, 1002.0371\]](#)

fit to PT:

$$a = \{0.05, 0.06, 0.08\} \text{ fm}$$

[PRELIMINARY]



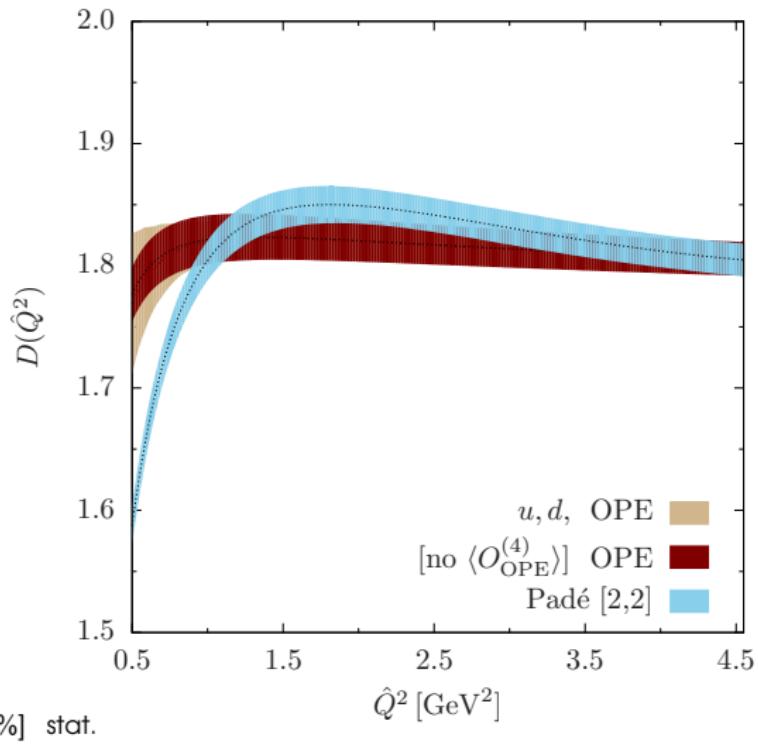
ongoing studies of systematic effects : lattice artefacts,  $Q^2$  interval,

order in OPE and perturbative expansions, ...

# fit to PT : comparison

$a \rightarrow 0$   
 $M_\pi^{\text{phys}}$

$u, d$



[PRELIMINARY]

# conclusions

- ▶ Adler function  $\leadsto a_\mu^{\text{had}}, \Delta\alpha_{\text{QED}}^{\text{had}}(Q^2), \alpha_s$
- ▶ good prospects for accurate determination of  $\Delta\alpha_{\text{QED}}^{\text{had}}(Q^2)$  on the lattice
- ▶ matching to perturbation theory : suffers both from statistical and systematic uncertainties

in view of future experimental results for  $a_\mu \dots$  or to address e.g. the  $e^+e^- - \tau$  difference  
→ improve precision and accuracy

- ▶ combination of standard and “mixed representation” methods  $\leadsto a_\mu^{\text{had}}$   
[talk by Anthony Francis]
- ▶ quark-disconnected diagrams [talk by Vera Gülpers]
- ▶ variance reduction techniques [poster by Eigo Shintani]
- ▶ ...



# Lattice VPF

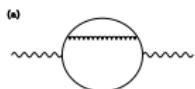
Local current

$$J_\mu^{(l, f)}(x) = Z_V \bar{\psi}_f(x) \gamma_\mu \psi_f(x)$$

conserved-local correlator

$$a^6 \left\langle \sum_{f=1}^{N_f} \left( Q_f J_\mu^{(ps, f)}(x) \right) \sum_{f'=1}^{N_f} \left( Q_{f'} J_\nu^{(l, f')}(0) \right) \right\rangle$$

$$\Pi_{\mu\nu}(\hat{Q}) = a^4 \sum_x e^{i\hat{Q}(x+a\hat{\mu}/2)} \langle J_\mu^{(ps)}(x) J_\nu^{(l)}(0) \rangle \quad \leadsto \quad \Pi(\hat{Q}^2)$$



$$\hat{Q}_\mu = \frac{2}{a} \sin \left( \frac{a Q_\mu}{2} \right)$$

# Adler function : combined fit

Adler function:

$$D(Q^2) = -12 \pi^2 Q^2 \frac{d\Pi(Q^2)}{dQ^2}$$

- **fit form :**

$$D(Q^2) = \text{Padé}(Q^2) [1 + \text{discr.} + \text{mass}] ,$$

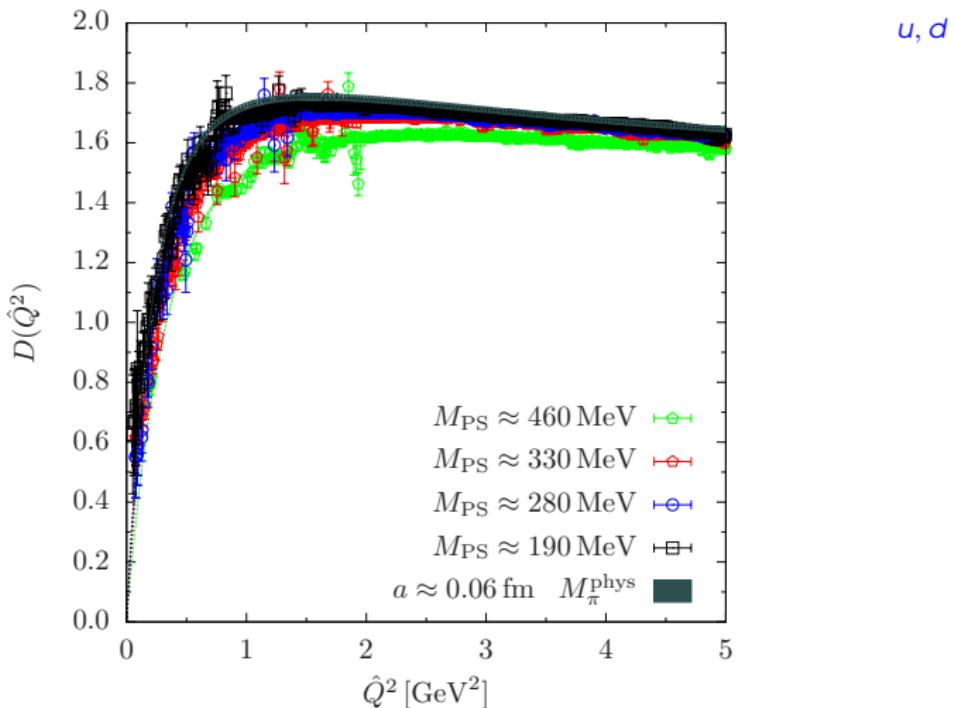
$$D(Q^2) = Q^2 \left( p_0 + \frac{p_1}{(p_2 + Q^2)^2} + \frac{p_3}{(p_4 + Q^2)^2} \right) \times \\ \left[ 1 + \left( d_1 a^n + d_2 (aQ)^n \right) + \left( \frac{c_1}{c_2 + Q^2} \right) (M_{\text{PS}}^2 - M_\pi^2) \right] .$$

$$n = \{1, 2\}$$

- consider 11 ensembles with different  $a, M_{\text{PS}}$
- $u, d, s_Q$  and  $c_Q$

# Adler function : light-quark mass dependence

$$D(Q^2) = \frac{3\pi}{\alpha} \frac{d}{d \log(q^2)} \Delta \alpha_{\text{QED}}^{\text{had}}(q^2)$$



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